Introduction to quantum computation with superconductors

FRANCESCO VISCHI NEST-CNR AND UNIVERSITY OF PISA

MAY 2017





A Q-bit is a quanto-mechanical two level system that can be manipulated in order to store and transport information. The two levels are indicated as |0> and |1>. The state is in general

$$|\psi\rangle = \cos\frac{\theta}{2}|0\rangle + e^{i\varphi}\sin\frac{\theta}{2}|0\rangle$$

Q-bits have been implemented with various physical systems: photon in cavity, NMR, quantum dots , etc. The most promising are trapped ions and superconducting circuits

➤A useful Q-bit must respect the DiVincenzo QC criteria, concerning scalability, initialization, long coherence times, measurement, manipulation



D. DiVincenzo

What's superconductivity

> It is a phase transition characterized by a critical temperature and zero resistance

> The electrons condense in a macroscopic quantum state. The condensate can be described by a complex order parameter $\psi(\vec{r})$, similar to a wave function.



Refs [Gr00,Ti04,Gr03]

Ideal JJ

One evidence of coherence is the Josephson effect in a Josephson Junction (JJ). Josephson junctions are devices where two superconducting parts are coupled by mean of a insulator (SIS)

The current and the voltage across a JJ is ruled by Josephson laws:

$$I = I_c \sin(\varphi)$$
$$V = \frac{\hbar}{2e} \frac{d\varphi}{dt}$$





Refs [Gr00,Ti04,Gr03]

CSJ model

The behavior of a JJ can be modeled better within the *RCSJ model* (Resistance-Capacitance Shunted Junction). We focus on zero noise and zero dissipation CSJ model.

$$\begin{split} H(N,\varphi) &= E_C N^2 - E_J \left(\cos(\varphi) + \frac{I}{I_C} \varphi \right) \\ \hline \text{Kinetic term} & \text{Tilted washboard potential} \end{split} \qquad E_C = \frac{(2e)^2}{2C} \qquad E_J = I_c \frac{\hbar}{2e} \end{split}$$



Tilted washboard potential



Refs [Ti04,Gr03]

Secondary macroscopic quantum effect

>In quantum mechanics, phase and number of particle are conjugate variables. It holds

$$\begin{split} [\hat{\varphi}, \hat{N}] &= i \\ \Delta \varphi \Delta N \geq 1 \end{split}$$

> The Hamiltonian can be quantized:

$$H(N,\varphi) = E_C \hat{N}^2 - E_J \left(\cos(\hat{\varphi}) + \frac{I}{I_C} \hat{\varphi} \right)$$

Refs [Ti04,Gr03]

Superconducting Q-bits

There are three types of superconducting Q-bits based on insulator JJ. In order of realization, with respective citation:

- Charge Q-bit P. Lafarge et al., Nat. 365, 422 (1993)
 Y. Nakamura et al. Nat. 398, 786 (1999))
- Flux Q-bit
 W. H. Caspar et al., Science 290, 773 (2000)
 J. R. Friedman et al., Nat. 406, 43 (2000)

Phase Q-bit J. M. Martinis et al., PRL 89, 117901 (2002)

FOR SIMPLICITY, IN THE FOLLOWING SLIDES WE'RE GOING TO DEAL WITH PRIMITIVE PROOF-OF-PRINCIPLE ARCHITECTURES. Many issues (e.g. coherence, noise, readout, coupling) have brought to more complex architectures

Phase Q-bit

The phase Q-bit is a single JJ driven by an external current. The junction state is brought near the voltage state, with $I \approx I_C$. The potential is then approximatively cubic with bound states

>Inizialization is accomplished by thermalization and adiabatic bias

> Manipulation of the state can be done by mean of the bias current.







Phase Q-bit: experimental facts

> Evidence of Rabi Oscillations depending on excitation pulse amplitude



Refs [Ma02]

The flux Q-bit is constituted by an insulated superconducting ring interrupted by JJ (simplest config).

► Working principle: flux quantization

$$\Phi_0 = \frac{h}{2e}$$

Let us consider for simplicity a single junction loop (alias RF-SQUID).. The Hamiltonian of the system is

$$H(N,\tilde{\Phi}|\Phi_e) = E_C N^2 + \frac{\tilde{\Phi}^2}{2L} - E_J \cos\left(\frac{2e}{\hbar}(\tilde{\Phi} - \Phi_e)\right)$$



Refs [Wa00, Fr00]

> The potential is a corrugated parabola. If the external flux is set to

$$\Phi_e = \frac{1}{2}\Phi_0 \qquad \Phi_0 = \frac{h}{2e}$$

the potential takes a quartic form (mexican hat).

Classically, the two fundamentals state correspond to current circulating clockwise or counterclockwise with degenerate energy.

The quantum tunneling removes degeneracy, so that the two lowest states are superpositions of $| \uparrow >$ and $| \downarrow >$.

>As function of the external parameter, the excitation energy shows a typical anticrossing behavior.





Refs [De04, We08]

Read out is made by a DC-SQUID that measures the threading flux if counterclockwise of clockwise

>Measurement of current switching under fixed tuned frequency returns a anti-crossing figure





Manipulation occurs by a coupled microwave line, that produces GHz oscillating magnetic field in the loop.

>Also in this case, Rabi oscillation were detected





Charge Q-bit

➢ In the charge Q-bit, the observable exploited for QC is the number of Cooper pairs in a superconducting small volume, called Cooper pair box (SCB)

> The Hamiltonian of a SCB under gating is

$$H(Q,\varphi) = E_C(N - N_G)^2 - E_J \cos(\varphi)$$





Charge Q-bit

The classical parabola branches are splitted by Josephson potential. In this region, QC can be realized

➢One read readout scheme consists in coupling a normal probe to the SCB with a high resistance tunnel junction. Then a −excedent- Cooper pair can decay in two quasiparticles and generate a detectable current in the probe





Refs [Na99, We08]

Charge Q-bit

With this readout scheme, coherent oscillations were observed with this measuring scheme [Na99]:

1. Polarize with a DC gate the SCB a little away from the anticrossing point

- 2. Give a sharp pulse in a pulse gate to make a non adiabatical transition
- 3. Keep the pulse voltage on for a desired time interval (so, let the system evolve)

4. Turn off the pulse. Now the system is in the desired superposition. Excited state will decay in two quasiparticles



Refs [Na99]

Summary

Superconductor devices are one way to implement QC, and the most promising after trapped ions

- The fundamental component is the JJ
- There are three types of Q-bits exploiting different form of hamiltonian (table)

Q-bit type	Potential polynomial approx	External param.	Parameter static component
Charge	Parabolic	Gate voltage	$V_G \approx \frac{E_C}{e}$
Flux	Quartic power (mexican hat)	Magnetic Flux	$\Phi \approx \frac{\Phi_0}{2}$
Phase	Third power	Current	$I \to I_C$



Thank you for the attention!

ANY QUESTION?

References

Superconductivity and Josephson effect

[Gr00] G. Grosso and G. Pastori Parravicini, Solid state physics, (Academic Press, 2000)

[Ti04] M. Tinkham, Introduction to superconductivity, (Dover, 2004)

[Gr03] R. Gross and A. Marx, Applied Superconductivity (http://www.wmi.badw.de/teaching/Lecturenotes/)

Evidence of quantization of phase and Cooper pair number [Ko82] R. H. Koch et al. , PRB 26, 75 (1982)
[Ko80] R. H. Koch et al. , PRL 45, 2132 (1980)
[Vo81] R. F. Voss et al. , PRL 47, 265 (1981)
[Ma87] J. M. Martinis et al. , PRB 35, 4682 (1987)

References

Quantum Q-bit, overviews and reviews

[Cl08] J. Clarke and F. K. Wilhelm, Nat. 453, 1031 (2008)
[We08] G. Wendin and V.S. Shumeiko, arXiv, arXiv:cond-mat/0508729v1
[De04] M. H. Devoret et al., arXiv, arXiv:cond-mat/0411174v1

Phase Q-bit

[Ma02] J. M. Martinis et al., PRL 89, 117901 (2002)

Flux Q-bit

[Wa00] C. H. van der Wal et al., Science 290, 773 (2000)
[Fr00] J. R. Friedman et al., Nature 406, 43 (2000)
[Ch03] I. Chiorescu et al., Science 299, 1869 (2003)

References

Charge Q-bit

[La93] P. Lafarge et al., Nat. **365**, 422 (1993) [Na99] Y. Nakamura et al. Nat. **398**, 786 (1999)