A new formulation of Lee-Wick models and its implications for quantum gravity



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Pre-thesis seminar

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Outline

- 1 Renormalizability and Quantum Gravity
- 2 Minkowski Higher-Derivative Theories
- **3** Lee-Wick Quantum Field Theory
- **4** New Formulation
- **5** Lee-Wick Models and Quantum Gravity
- 6 Conclusions and Future Developments



$$\begin{array}{c} \text{QED} \\ A_{\mu} \\ F_{\mu\nu} = \partial_{\mu}A_{\nu} - \partial_{\nu}A_{\mu} \\ \mathcal{L} = -\frac{1}{4}F_{\mu\nu}^{2} + \bar{\psi}(i\partial \!\!\!/ - eA \!\!\!/ - m)\psi, \quad \alpha = \frac{e^{2}}{4\pi} \end{array}$$

$$\mathcal{L}_R = -\frac{1}{4}F'^2_{\mu\nu} + \bar{\psi}'(i\partial \!\!\!/ - e'A\!\!\!/ - m')\psi',$$

QED

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$$g_{\mu\nu}$$

$$R^{\mu}_{\ \nu\rho\sigma} \simeq \partial\partial g + \partial g\partial g + \dots$$

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$$\begin{array}{ccc} \text{QED} & \text{Quantum Gravity} \\ A_{\mu} & g_{\mu\nu} \\ F_{\mu\nu} = \partial_{\mu}A_{\nu} - \partial_{\nu}A_{\mu} & R^{\mu}_{\ \nu\rho\sigma} \simeq \partial\partial g + \partial g\partial g + \dots \\ \mathcal{L} = -\frac{1}{4}F^{2}_{\mu\nu} + \bar{\psi}(i\partial \!\!\!/ - eA \!\!\!/ - m)\psi, \quad \alpha = \frac{e^{2}}{4\pi} & \mathcal{L} = -\frac{1}{2\kappa^{2}}\sqrt{-g}R, \quad \kappa^{2} = 8\pi G \\ \mathcal{L}_{R} = -\frac{1}{4}F^{\prime 2}_{\mu\nu} + \bar{\psi}'(i\partial \!\!\!/ - e'A' - m')\psi', & \mathcal{L}_{R} = -\frac{1}{2\kappa^{2}}\sqrt{-g}[R + c_{1}R^{2} + c_{2}R_{\mu\nu}R^{\mu\nu} \\ + c_{3}R^{\mu\nu}_{\rho\sigma}R^{\alpha\beta}_{\alpha\beta}R^{\alpha\beta}_{\mu\nu} + \underbrace{\cdots}_{\infty}]. \end{array}$$

In QG renormalization generates an infinite number of new counterterms. A high-energy modification of the theory is necessary.

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- Higher derivatives \Rightarrow improve the convergence;
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Minkowski higher-derivative theories

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U.G. Aglietti and D. Anselmi, *Inconsistency of Minkowski higher-derivative theories*, Eur. Phys. J. C 77 (2017) 84, 16A2 Renormalization.com and arXiv:1612.06510 [hep-th]. Inconsistencies in Minkowski formulation:

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The theory cannot be defined directly on Minkowski spacetime.

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Our formulation solves all these problems.

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Modify the contour integration in loop integrals (T.D. Lee and G.C. Wick).

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$$iD(p^2, m^2, \epsilon) = \frac{iM^4}{(p^2 - m^2 + i\epsilon)((p^2)^2 + M^4)}$$

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II) Pinching singularities



$$\mathcal{J}(p) = \int \frac{\mathrm{d}^D k}{(2\pi)^D} D(k^2, m_1^2, \epsilon_1) D((k-p)^2, m_2^2, \epsilon_2)$$
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II) Pinching singularities

CLOP prescription (Cutkosky at al.): two different scales M and M' s.t.

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Solution: define LW models as Nonanalytically Wick rotated Euclidean theories.

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New formulation

Analytic structure of $\mathcal{J}(p)$ in the p^0 complex plane

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Analytic structure of $\mathcal{J}(p)$ in the p^0 complex plane $\mathbf{p} = 0$ $\lim [p^0]$ $\sqrt{2}M$ $\operatorname{Re}[p^0]$ $\sqrt{2}M$

The amplitude is ill-defined on the real axis above the threshold $p^2 = 2M^2$.

New formulation

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Analytic structure of $\mathcal{J}(p)$ in the p^0 complex plane

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Lorentz invariance seems violated (already noticed by Nakanishi) Deform the branch cuts \leftrightarrow Deform the k integration domain



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Constraints on the deformation:

• Symmetric w.r.t. the real axis (unitarity).

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The amplitude are well defined but nonanalytic.

The deformation is practically hard to implement

We argue that

$$\mathcal{J}_{\mathrm{LW}}^{>}(p) = \frac{1}{2} \left[\mathcal{J}_{\mathrm{LW}}^{0+}(p) + \mathcal{J}_{\mathrm{LW}}^{0-}(p) \right]$$

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No prescription with nonvanishing a is consistent with this formulation. CLOP gives $a=\pi {\rm sgn}(M'-M)$

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Discrepancy above the threshold



Our formulation gives physical predictions which differ from the previous ones.

$$\mathcal{L}_{\rm QG} = -\frac{1}{2\kappa^2} \sqrt{-g} \Big[R - \frac{1}{M^4} \big(D_\rho R_{\mu\nu} \big) \big(D^\rho R^{\mu\nu} \big) + \frac{1}{2M^4} \big(D_\rho R \big) \big(D^\rho R \big) \Big]$$

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Expansion $g_{\mu\nu} = \eta_{\mu\nu} + 2\kappa h_{\mu\nu}, \qquad \eta_{\mu\nu} = \text{diag}(1, -1, -1, -1).$

Propagator in harmonic gauge

$$\langle h_{\mu\nu}(p)h_{\rho\sigma}(-p)\rangle^{\text{free}} = \frac{iM^4}{2(p^2+i\epsilon)} \frac{\eta_{\mu\rho}\eta_{\nu\sigma} + \eta_{\mu\sigma}\eta_{\nu\rho} - \eta_{\mu\nu}\eta_{\rho\sigma}}{(p^2)^2 + M^4}.$$

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$$\mathcal{L}'_{\rm QG} = -\frac{1}{2\kappa^2} \Big[2\Lambda_C + \zeta R - \frac{1}{M^2} R_{\mu\nu} P_n (\Box_c/M^2) R^{\mu\nu} \\ + \frac{1}{2M^2} R Q_n (\Box_c/M^2) R + \mathcal{V}(R, M, \alpha_i) \Big]$$

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$$\mathcal{L}_{\rm QG} = -\frac{1}{2\kappa^2} \sqrt{-g} \Big[R - \frac{1}{M^4} \big(D_{\rho} R_{\mu\nu} \big) \big(D^{\rho} R^{\mu\nu} \big) + \frac{1}{2M^4} \big(D_{\rho} R \big) \big(D^{\rho} R \big) \Big]$$

Expansion $g_{\mu\nu} = \eta_{\mu\nu} + 2\kappa h_{\mu\nu}, \qquad \eta_{\mu\nu} = \text{diag}(1, -1, -1, -1).$

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Problem of uniquness

Superrenormalizable models of quantum gravity are infinitely many.

New quantization prescription

D. Anselmi, On the quantum field theory of gravitational interactions, JHEP 1706 (2017) 086, 17A3 Renormalization.com and arXiv:1704.07728 [hep-th].

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$$\mathcal{L} = \frac{1}{2} \partial_{\mu} \varphi \partial^{\mu} \varphi - \frac{\lambda}{4!} \varphi^4$$

Modified Euclidean propagator

$$\frac{p_E^2}{(p_E^2)^2 + \mathcal{E}^4}, \quad \mathcal{E} = \text{ ficticious LW scale}$$

New prescription

$$\lim_{\mathcal{E}\to 0} \frac{p^2}{[(p^2)^2 + \mathcal{E}^4]_{\rm LW}}$$

$$\mathcal{L}_{QG} = -\frac{1}{2\kappa^2} \sqrt{-g} \Big[2\Lambda_C + \zeta R - \frac{\gamma}{M^2} R_{\mu\nu} R^{\mu\nu} + \frac{1}{2M^2} (\gamma - \eta) R^2 \Big], \quad \zeta > 0, \ \gamma < 0.$$

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$$\left\{\frac{1}{p^2+i\epsilon}-\frac{\gamma(\zeta M^2+\gamma p^2)}{\left[(\zeta M^2+\gamma p^2)^2+\mathcal{E}^4\right]_{\rm LW}}\right\}\frac{i}{2\zeta}(\eta_{\mu\rho}\eta_{\nu\sigma}+\eta_{\mu\sigma}\eta_{\nu\rho}-\eta_{\mu\nu}\eta_{\rho\sigma}).$$

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A unique, renormalizable and unitary theory of QG in 4 dim.

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A unique, renormalizable and unitary theory of QG in 4 dim.

Unitarity can be proved only if $\Lambda_C = 0$ but realistic models have $\Lambda_C \neq 0$. The cosmological constant might be an anomaly of unitarity.

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Conclusions and future developments

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- Study the physics of the new model of QG which involves our formulation (in preparation)
- Investigate the new possible phenomenology of fundamental interactions due to the new quantization prescription.