The Fuzzball conjecture

Lorenzo Bartolini

The Information Paradox

Building Black Holes in String Theory

The Fuzzball Conjecture

Conclusions

The Black Holes' Information Paradox and String Theory Based on works by S.Mathur, O.Lunin, K.Skenderis, I.Bena, S.Giusto, J.F.Morales, M.Bianchi

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21 September 2017

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The simplest BH anatomy

$$ds^{2} = \left(1 - \frac{R_{s}}{r}\right)c^{2}dt^{2} - \left(1 - \frac{R_{s}}{r}\right)^{-1}dr^{2} - r^{2}d\Omega_{2}^{2}$$

• Completely determined by
$$R_s \equiv \frac{2MG}{c^2}$$

• Real singularity at r = 0

More general cases completely determined by M, Q, J

 \mathbb{T}

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Throw a box of hot gas inside the BH $\downarrow\downarrow$

It had a definite entropy, now it seems lost

Save the II Law of Thermodynamics: assign an entropy to Black Holes

$$\boxed{ \begin{array}{c} \hline \text{Bekenstein 1973} \end{array} } \\ S_{Bek} = \frac{A}{4G} \Rightarrow \frac{dS_{Bek}}{dt} + \frac{dS_{matter}}{dt} \geq 0 \end{array} } \end{array}$$

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Entropy puzzle

Puzzle: entropy has usually a fundamental interpretation



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There should be
$$\mathcal{N} = \exp{S} \simeq \exp{\left[10^{76} \left(\frac{M}{M_{\odot}}\right)^2\right]}$$

microstates

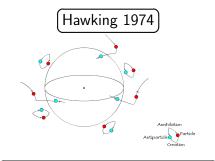
However...

(BLACK HOLES HAVE NO HAIR)

No real trouble: "hair" may be produced by Quantum Gravity and hidden well beyond the horizon

Hawking Radiation

Semiclassically, Black Holes are no longer stable



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QFT+gently curved background

Pair creation near the horizon: one falls in the BH, the other escapes to infinity

Seen by a far observer as thermal emission from the BH

Black Hole Thermodynamics

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Law	Thermodynamics	Black Holes
Zeroth	T constant over bodies	κ constant over
	at equilibrium	the horizon
First	dE = TdS - PdV	$dM = \frac{1}{8\pi}\kappa dA + \Omega dJ$
Second	$\delta S \leq$ 0	$\delta A \leq 0$
Third	T = 0 not achievable	$\kappa = 0$ not achievable

To complete the duality, we have

Hawking emission \leftrightarrow Blackbody radiation

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Throw some matter in a BH, wait for complete evaporation through Hawking emission

- Initial matter ightarrow pure state $|\psi
 angle$
- Final state \rightarrow density matrix ρ_{rad}

The final state carries no information of what fell beyond the horizon

(Loss of UNITARY EVOLUTION)

Based on very few, reasonable assumptions

- Quantum Gravity effects are confined to lp
- Vacuum is unique

Details of Quantum Gravity are not relevant

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How can Quantum Gravity work around the paradox?

- \blacktriangleright Planck sized remnants \rightarrow ugly: size independent from S
- ▶ Non-local interactions (singularity ↔ horizon)
- Unitarity is effectively lost
- Structure modified up to the horizon

String theory points towards the last one

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String Theory in few words

One principle: fundamental objects are 1-dimensional

Consequences

- SUSY to produce fermions
- Spacetime dimension= $10 \Rightarrow$ compact dimensions
- p-Branes: extended, charged objects on which strings can end
- In principle: no dimensionless free parameters
- Five theories related by web of dualities \Rightarrow M-Theory?

But the most important

It is a consistent Quantum Gravity (SUGRA as Eff.Th.)

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Conclusions

it MUST solve the information paradox somehow

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Building Black Holes (Susskind)

How to solve the entropy puzzle:

Take a system of fundamental objects (strings/branes)

- Count degeneracy \rightarrow Microscopic count: easy at small g_s
- ▶ Increase g_s until a BH is formed→ Compute Area of Horizon
- Compare the degeneracy with the Bekenstein entropy: should match (or we throw away string theory)

Inconsistency: levels can shift with changes in g_s \downarrow Focus on BPS states: levels shift together (Thank y

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Building Black Holes (Susskind)

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2-Charge state (winding n_y , momentum n_p in N_L)

$$S_{micro} = 2\sqrt{2}\pi\sqrt{n_y n_p}$$

$$ds_{naive}^{2} = \left(1 + \frac{Q_{1}}{r^{2}}\right)^{-1} \left[-dudv + \frac{Q_{p}}{r^{2}}dv^{2}\right] + dx_{i}dx^{i} + dz_{a}dz^{a}$$

No horizon: if we add higher derivative terms, it develops a horizon at r = 0 with

$$S_{Bek} = S_{micro}$$

We called it naive for a fundamental reason

This metric cannot exist in string theory

(datails later, crucial for next section)

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3-Charge state (add
$$n_5$$
 5-Branes)

 $S_{micro} = 2\pi \sqrt{n_y n_p n_5}$

Construct the corresponding metric, compute the horizon area

 $S_{Bek} = S_{micro}$

No subtleties, remarkable matching and moreover...

This solution is an extremal BH \rightarrow no Hawking radiation Perturb this configuration to make it not extremal (add some N_R)

Left and Right moving excitations can collide and escape as massless quanta leaving the Brane configuration

Rate and spin dependence agree with Hawking emission

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Real geometries

$$ds_{naive}^{2} = \left(1 + \frac{Q_{1}}{r^{2}}\right)^{-1} \left[-dudv + \frac{Q_{p}}{r^{2}}dv^{2}\right] + dx_{i}dx^{i} + dz_{a}dz^{a}$$

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Cannot be produced by any string configuration

Real geometries

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Cannot be produced by any string configuration (WHY?)

 ds_{naive}^2 is spherically symmetric \Rightarrow generated by a pointlike source in transverse space

But the momentum MUST be carried by transverse waves on the string

 \Downarrow

Cannot be pointlike in transverse space Should be considered just a valid SUGRA solution far from r = 0 The Fuzzball conjecture

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Vibration profile described by $\vec{F}(v \equiv t - y)$ with $0 \le y \le 2\pi n_y R$

- Find the corresponding metric
- Make approximations (neighbouring "strands" give the same contribution)
- ► Use dualities of string theory to move to D1 D5 system

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Real geometries

Resulting solution

$$ds^{2} = \sqrt{\frac{H}{1+K}} \left[-(dt - A_{i} dx^{i})^{2} + (dy + B_{i} dx^{i})^{2} \right] + \sqrt{\frac{1+K}{H}} dx_{i} dx^{i} + \sqrt{H(1+K)} dz_{a} dz^{a}$$

$$\begin{aligned} H^{-1} &= 1 + \frac{\mu Q_1}{L} \int_0^{\mu L} \frac{dv}{|\vec{x} - \mu \vec{F}(v)|^2} \\ \mathcal{K} &= \frac{\mu Q_1}{L} \int_0^{\mu L} \frac{dv (\mu^2 \dot{F}_i(v))^2}{|\vec{x} - \mu \vec{F}(v)|^2} \\ \mathcal{A}_i &= \frac{\mu Q_1}{L} \int_0^{\mu L} \frac{dv \mu \dot{F}_i}{|\vec{x} - \mu \vec{F}(v)|^2} \quad ; \quad dB = - \star_4 dA \end{aligned}$$

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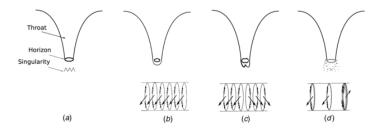
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How do they look?



The geometry can be divided in three main regions

- $r \to \infty \Rightarrow$ FLAT SPACE
- ▶ Intermediate "throat" $AdS_3 \times S_3 \times T_4 \Rightarrow \approx ds_{naive}^2$
- $r \rightarrow 0 \Rightarrow$ ends in a smooth "cap"

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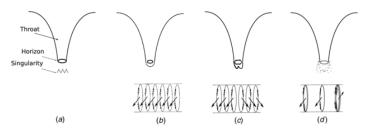
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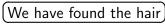
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How do they look?





• Different $\vec{F}(v) \Rightarrow$ Different caps

- From afar, they look like the classical "naive" BH
- Horizonless \Rightarrow Crucial: microstates have no entropy

The classical BH is a coarse-grained description

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The size

Last issue

Information of the microstates must reach the horizon scale, or we have not solved the paradox

$$\Delta x = |\dot{\vec{F}}| \Delta y \sim \sqrt{\alpha'} \equiv \ell_s$$

For
$$x \gg \sqrt{lpha'} \Rightarrow ds^2 o ds^2_{naive}$$

Compute area at $|\vec{x}| = \sqrt{\alpha'}$ in transverse space

$$rac{A}{4G_{10}}\sim \sqrt{n_y n_p}=S_{micro}$$

Area reproduces the Bekenstein Entropy!

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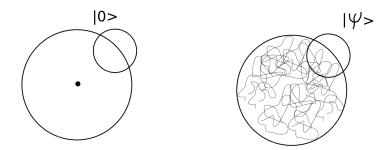
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The size



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Conclusions

The state near the horizon is not the vacuum, instead we have

$$\langle 0|\psi
angle pprox 0$$

The "cap" region has a size which satisfies a Bekenstein relation

Does the boundary of this region behave like a horizon?

Simple test: consider a quantum falling in the cap region

The time it takes to cross the generic cap is found to be proportional to $\left(\frac{M}{m_{pl}}\right)^2$

A classical observer for which $\hbar \sim$ 0 finds an infinite crossing time

Everything that falls in the cap, can never escape

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String theory passes a nontrivial test as a consistent Quantum Gravity

- BHs are coarse-grained descriptions of Fuzzballs
- Entropy comes from coarse-graining of N = e^S
 Fuzzball states
- Fuzzball states are horizonless (no information loss)
- ► Effects from string theory extend to horizon scale ⇒ the "would-be horizon" carries hair
- Hawking radiation is no different from that emitted by a burning coal

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Thanks for your attention