





Experimental studies of the lifetimes of Rydberg atoms in cold atomic gases

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General introduction



Spontaneous decay.

General introduction



Rydberg atoms

Exaggerated properties with respect to ground state atoms.

Scaling with $n^* = (n - \delta_{qdt})$:

- Binding energy: (n*)⁻²
- Orbital radius: (n*)²
- Lifetime: (n*)³
- Polarizability: (n*)⁷
- Van der Waals coefficient: (n*)¹¹

Several applications:

- Many body physics;
- Quantum technologies.



Two effects contribute to the lifetime of Rydberg atoms:

Spontaneous emission



Blackbody induced transitions



ground state

Two effects contribute to the lifetime of Rydberg atoms:

Spontaneous emission





$$\frac{1}{\tau_{eff}} = \Gamma_0 + \Gamma_{BBR} = \frac{1}{\tau_0} + \frac{1}{\tau_{BBR}}$$



Two effects contribute to the lifetime of Rydberg atoms:



The experimental apparatus

Preparation: Atoms cooled and trapped with lasers and magnatic fields: magneto-optical trap





Excitation: Excited with two photon excitation from ground state Rydberg state

$$\Omega_{tot} = \sqrt{\left(\frac{\Omega_{421}\Omega_{1012}}{2\Delta_{6P}}\right)^2 + \Delta^2}$$

The experimental setup

Detection: Rydberg atoms are field ionized and accelerated towards a channeltron



The channeltron converts an incident ion to an electrical signal



The experimental setup

Detection: Rydberg atoms are field ionized and accelerated towards a channeltron



We can record:

- Mean number of ions <N>;
- Standard deviation σ;
- Mandel paramenter $Q = \frac{\sigma^2}{\langle N \rangle} 1$;
- Arrival time distribution.



The experimental setup

Two main parts for detection:

- Ionization of Rydberg atoms;
- Acceleration to channeltron.





First we focus on the trajectories.

lons produced by photoionization from 5P











Characterization of the detection: electric fields in the cell

Electric fields in the cell:

- Complex configuration of potentials inside the apparatus;
- Quartz cell is dielectric, screening effects.

We measure the Stark shift to know the actual field on atoms: $1 - \frac{1}{2}$

$$\Delta E = -\frac{1}{2}\alpha \mathcal{E}^2$$

Characterization of the detection: electric fields in the cell

Electric fields in the cell:

- Complex configuration of potentials inside the apparatus;
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We measure the Stark shift to know the actual field on atoms: $1 - 2^{2}$

$$\Delta E = -\frac{1}{2}\alpha \mathcal{E}^2$$

Loss of atom from the MOT:

$$\frac{dN(t)}{dt} = L - \Gamma_0 N(t) - \Gamma_{Ryd} N(t)$$
$$\Gamma_{ryd}(\nu) = \Gamma_0 \left(\frac{N_0}{N_{stat}(\nu)} - 1\right)$$



Characterization of the detection: electric fields in the cell



Fit of the data on the quadratic Stark shift:

$$\Delta E = -\frac{1}{2}\alpha \mathcal{E}^2$$

$$f(\nu) = y_0 - \frac{1}{2}\alpha_{70S} \left(\gamma V - E_0\right)^2$$

$$\alpha_{70S} = 557.4 \frac{MHz}{Vcm^{-2}}$$

 $E_0 = (-6.7 \pm 0.2) \times 10^{-2} \,\mathrm{V/cm}$ $\gamma = (4.52 \pm 0.02) \times 10^{-3} \,\mathrm{cm}^{-1}$

Characterization of the detection: ionization threshold



We count the number of atoms in a Rydberg state as a function of time, and obtain an exponential decay.

Typical measurement: Switch on the ionization fields with slowly increasing sweep. Rydberg atoms ionized in different moments.



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Difficult to apply for high n because ionization thresholds are very near.

We use the *depump technique*

The lifetime of Rydberg atoms: depump technique



Measurement of the lifetime

Two parts:

- Excite atoms, wait and ionize all Rydberg atoms;
- Excite atoms, wait, depump atoms in the <u>target state</u>, and ionize the other (<u>support states</u>)



Measurement of the lifetime: lifetime of the target state



Measurement of the lifetime: lifetime of the target state+support states



Measurement of the lifetime: effects of interactions



Dependence of the target+support lifetimes on the number of initial excitations



Anomalous initial fast decay, deviation from simple exponential decay for high lying Rydberg states.

Conclusions

- I have presented the results on the measurements of the lifetimes of Rydberg atoms;
- Good agreement with theoretical calculation between n=60 and n=80;
- Unexpected deviation from the theory for n>80;
- Dependence with number of Rydberg atoms and non-exponential decay suggest effects of interactions.

The results obtained during my PhD led to the following publications:

- De-excitation spectroscopy of strongly interacting Rydberg gases, C. Simonelli, M. Archimi, L. Asteria, D. Capecchi, G. Masella, E. Arimondo, D. Ciampini, O. Morsch, arXiv:1707.01382, submitted to Phys. Rev. A (2017);
- Experimental signatures of an absorbing-state phase transition in an open driven many-body quantum system, R. Gutierrez, C. Simonelli, M. Archimi, F. Castellucci, E. Arimondo, D. Ciampini, M. Marcuzzi, I. Lesanovsky, O. Morsch, arXiv:1611.03288, submitted to Phys. Rev. A (2017).